

The Laplace Transform: Calculus

Jonathan L.F. King
 University of Florida, Gainesville FL 32611-2082, USA
 squash@ufl.edu
 Webpage <http://squash.1gainesville.com/>
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ABSTRACT: This first gives a defn of *exponential order* which is better adapted to convolution. Following, is a discussion of the “tapping on a bell” problem; one text called this “soldiers marching on a bridge”. Both interpretations need a grain-of-salt...

Prelims. On a (possibly infinite) interval $J \subset \mathbb{R}$, a function $f: J \rightarrow \mathbb{C}$ is *locally-integrable* if, for each bounded subinterval $[a, b] \subset J$, integral $\int_a^b f$ exists and is finite. A *sufficient* (but not necessary) condition is that on each $[a, b]$, our f is bounded with only finitely-many discontinuities.

For a fnc $f: [0, \infty) \rightarrow \mathbb{C}$ and complex number s we define the “Laplace transform of f ”, evaluated at s ,

$$1: \quad \hat{f}(s) = [\mathcal{L}(f)](s) := \int_0^\infty e^{-st} \cdot f(t) \cdot dt,$$

for those values of s where this integral exists.

For complex numbers α and β , let $\alpha \succ \beta$ mean $\operatorname{Re}(\alpha) > \operatorname{Re}(\beta)$. For a real μ , say that “ f has *exponential order μ* ”, written $\boxed{f \in \operatorname{Ord}(\mu)}$, if $f: [0, \infty) \rightarrow \mathbb{C}$ and f is locally-integrable and

$$\dagger: \quad \forall Q > \mu: \quad \lim_{t \rightarrow \infty} |f(t)|/e^{Qt} = 0.$$

One can replace \dagger by the seemingly weaker

$$\ddagger: \quad \forall Q > \mu: \quad \limsup_{t \rightarrow \infty} |f(t)|/e^{Qt} < \infty,$$

but they are [exercise] equivalent.

2: Lemma. Consider an $f \in \operatorname{Ord}(\mu)$. Then $\hat{f}(s)$ exists for each s with $\operatorname{Re}(s) > \mu$. Indeed the integrand in (1) is absolutely integrable. \diamond

Proof. Fix s with $x := \operatorname{Re}(s) > \mu$, then pick Q with $x > Q > \mu$. Our t is real, so $|e^{-st}| = e^{-xt}$. Hence the integrand in (1) is eventually bounded,

$$\left| e^{-st} \cdot f(t) \right| < e^{-xt} \cdot e^{Qt} \stackrel{\text{note}}{=} e^{-[x-Q]t},$$

once t is large enough. Since $x-Q$ is positive, this last is integrable over $t \in [0, \infty)$. \diamond

Convolution. Recall that the (one-sided) *convolution* of two (locally-integrable) functions $f, g: [0, \infty) \rightarrow \mathbb{C}$, is the function

$$3a: \quad [f \otimes g](t) := \int_0^t f(t-v) \cdot g(v) \, dv.$$

3b: Lemma. Suppose $f, g \in \operatorname{Ord}(\mu)$. Then the convolution $f \otimes g \in \operatorname{Ord}(\mu)$. \diamond

[Proof omitted.]

We need the following standard tool, to prove a uniqueness result.

4a: Identically-zero lemma. On a closed bounded interval $J := [b, c] \subset \mathbb{R}$, consider *continuous* functions $h, G: J \rightarrow \mathbb{C}$.

i: Suppose $h \geq 0$. If $\int_J h = 0$, then $h \equiv 0$, i.e, h is *identically-zero*. [Exer:]

ii: If G is real-valued and $\int_J [G \cdot G] = 0$, then $G \equiv 0$. [Exer:]

iii: Suppose

$$4b: \quad \forall n \in \mathbb{N}: \quad \int_J x^n \cdot G(x) \, dx = 0.$$

Then G is identically-zero on J . \diamond

Proof of (iii). Splitting G into real and imaginary parts, WLOG G is real-valued. Stmt (4b) implies, for each polynomial p , that $\int_J [p \cdot G]$ is zero.

By the Weierstrass Approximation Thm, there is a sequence of (real) polynomials, so that $p_k \xrightarrow{k \rightarrow \infty} G$ *uniformly*. Consequently,

$$\int_J [G \cdot G] \stackrel{\text{Exer.}}{=} \lim_{k \rightarrow \infty} \int_J [p_k \cdot G] = \lim_{k \rightarrow \infty} 0 = 0. \quad \blacklozenge$$

4c: Uniqueness Thm. Consider fncs $g, \gamma \in \operatorname{Ord}(\mu)$. Suppose their Laplace transforms, \hat{g} and $\hat{\gamma}$, agree on some real interval $\mathbf{I} := [s_0, \infty)$. If g and γ are continuous, then $g = \gamma$. \diamond

Proof. The difference fnc $f := g - \gamma$ is cts and in $\text{Ord}(\boldsymbol{\mu})$, with \widehat{f} identically-zero on \mathbf{I} . So our goal is to show f identically-zero on $[0, \infty)$.

WLOG $s_0 > \boldsymbol{\mu}$. Define $G(u) := u^{s_0} \cdot f(-\log(u))$, and $G(0) := 0$. CoV $t = -\log(u)$ shows that

$$\lim_{u \searrow 0} G(u) = \lim_{t \nearrow \infty} e^{-s_0 t} f(t)$$

equals 0, since $s_0 > \boldsymbol{\mu}$. Hence $\boxed{G \text{ is cts on } [0, 1]}$.

CoV $t = -\log(u)$ makes $dt = -du/u$ and

$$e^{-st} = [e^{-t}]^s = u^s.$$

Evaluating at $s+1$, the CoV rewrites (1) as

$$\begin{aligned} \widehat{f}(s+1) &= \int_1^0 u^{s+1} \cdot f(-\log(u)) \cdot \frac{-du}{u} \\ &= \int_0^1 u^s \cdot f(-\log(u)) du \\ &= \int_0^1 u^{s-s_0} \cdot G(u) du. \end{aligned}$$

For $n = 0, 1, 2, \dots$, let $s := s_0 + n$ and note $s+1$ is in \mathbf{I} . Hence $\widehat{f}(s+1)$ is zero, i.e. $\int_0^1 u^n G(u) du$ is zero. By (4a), then, $G \equiv 0$ on $[0, 1]$.

So on half-open $(0, 1]$, map $[u \mapsto f(-\log(u))]$ is zero. Hence $f \equiv 0$ on $[0, \infty)$. \blacklozenge

5a: TL-Diff. Fix a natnum N . Taking the N^{th} derivative,

$$\begin{aligned} \left[\frac{d}{ds} \right]^N \widehat{f}(s) &= \mathcal{L}([t]^N \cdot f(t))(s) \\ &\stackrel{\text{note}}{=} [-1]^N \cdot \mathcal{L}(t^N \cdot f(t))(s). \quad \blacklozenge \end{aligned}$$

5b: LT-Diff. Suppose f differentiable. Then

$$\widehat{f}'(s) = [s \cdot \widehat{f}(s)] - f(0).$$

For $N \in \mathbb{N}$, suppose that an N -times differentiable f , has $f, f', \dots, f^{(N-1)} \in \text{Ord}(\boldsymbol{\mu})$. Then

$$\widehat{f^{(N)}}(s) = [s^N \cdot \widehat{f}(s)] - \sum_{\substack{j+k=N-1 \\ j,k \in \mathbb{N}}} s^j \cdot [f^{(k)}(0)],$$

where the sum is over all ordered-pairs (j, k) of natnums. \blacklozenge

5c: LT-Exp. Fix $B \in \mathbb{C}$. Then

$$\mathcal{L}(e^{Bt} \cdot f(t))(s) = \widehat{f}(s - B),$$

whenever $s \succ B + \boldsymbol{\mu}$. \blacklozenge

5d: LT-Convolve. $[\mathcal{L}(f \circledast g)](s) = \widehat{f}(s) \cdot \widehat{g}(s)$. \blacklozenge

5e: TL-Integrate. For real s with $s > \boldsymbol{\mu}$,

$$\int_s^\infty \widehat{f}(u) du = \mathcal{L}\left(\frac{f(t)}{t}\right)(s) \quad \blacklozenge$$

Proof of (5e). For t positive, note

$$\begin{aligned} * : \int_s^\infty e^{-ut} du &= \frac{1}{t} \cdot e^{-ut} \Big|_{u=\infty}^{u=s} \\ &= \frac{1}{t} \cdot [e^{-st} - e^{-\infty}] = \frac{1}{t} \cdot e^{-st}. \end{aligned}$$

Applying the definition, $\int_s^\infty \widehat{f}(u) du$ equals

$$\begin{aligned} \int_s^\infty \left[\int_0^\infty e^{-ut} \cdot f(t) \cdot dt \right] du &= \int_0^\infty f(t) \int_s^\infty e^{-ut} du dt \\ &\stackrel{\text{by } (*)}{=} \int_0^\infty f(t) \cdot \frac{1}{t} \cdot e^{-st} dt. \end{aligned}$$

And this is the defn of $\mathcal{L}\left(\frac{f(t)}{t}\right)(s)$. \blacklozenge

Pf of (5a), TL-Diff. With appropriate conditions on f , we can differentiate under the integral sign in (1), applying $\frac{d}{ds}$, to get the $N=1$ case of (5a). Now induct on N . \blacklozenge

Pf LT-Diff. Our $\widehat{f}'(s)$ is the limit, as $M \nearrow \infty$, of

$$\int_0^M \underbrace{e^{-st}}_u \cdot \underbrace{f'(t)}_{dv} dt.$$

Integrating by parts, this equals

$$\begin{aligned} * : & \underbrace{e^{-st}}_u \cdot \underbrace{f(t)}_v \Big|_{t=0}^{t=M} - \int_0^M \underbrace{f(t)}_v \cdot \underbrace{[-s]e^{-st}}_{du} dt \\ &= \left[s \cdot \int_0^M f(t) \cdot e^{-st} dt \right] - \frac{f(t)}{e^{st}} \Big|_{t=M}^{t=0}. \end{aligned}$$

This last term equals $\frac{f(0)}{1} - \frac{f(M)}{e^{sM}} = f(0) - \frac{f(M)}{e^{sM}}$. If $\text{Re}(s) > \mu$, then $\frac{f(M)}{e^{sM}} \rightarrow 0$ as $M \nearrow \infty$. Sending $M \nearrow \infty$ thus sends $(*)$ to $[s \cdot \widehat{f}(s)] - f(0)$, as desired.

Finally, the formula for $\widehat{f^{(N)}}$ follows by induction on N . ♦

Pf LT-Exp. By defn, $\mathcal{L}(e^{Bt} \cdot f(t))(s)$ equals

$$\int_0^\infty e^{-st} e^{Bt} f(t) dt = \int_0^\infty e^{-(s-B)t} f(t) dt.$$

This last integral converges once $\text{Re}(s - B) > \mu$, i.e, once $s \succ B + \mu$. ♦

Proof of LT-Convolve. Define $\mathbf{1}(\text{true}) := 1$ and $\mathbf{1}(\text{false}) := 0$. We can now write RhS(3a) as

$$\int_0^\infty \mathbf{1}(v \leq t) \cdot f(t-v) \cdot g(v) dv.$$

Hence $[\mathcal{L}(f \otimes g)](s)$ equals

$$\int_0^\infty e^{-st} \int_0^\infty \mathbf{1}(v \leq t) \cdot f(t-v) \cdot g(v) dv dt.$$

Under mild conditions on f and g , we can reverse the integrals, giving that $[\mathcal{L}(f \otimes g)](s)$ equals

$$\int_0^\infty g(v) \int_0^\infty \mathbf{1}(v \leq t) \cdot e^{-st} \cdot f(t-v) dt dv.$$

The inner integral can be written as

$$\int_v^\infty e^{-st} \cdot f(t-v) dt.$$

The change of variable $x = t - v$ gives that $dx = dt$. So this integral becomes

$$\begin{aligned} \int_0^\infty e^{-s[x+v]} \cdot f(x) dx &= e^{-sv} \int_0^\infty e^{-sx} \cdot f(x) dx \\ &= e^{-sv} \cdot \widehat{f}(s). \end{aligned}$$

Thus $[\mathcal{L}(f \otimes g)](s)$ equals

$$\int_0^\infty g(v) \cdot e^{-sv} \cdot \widehat{f}(s) dv \stackrel{\text{note}}{=} \widehat{f}(s) \cdot \widehat{g}(s). \quad \blacklozenge$$

Examples

Below, $B \in \mathbb{C}$ and $F \in \mathbb{R}$. Easily $[\mathcal{L}(1)](s)$ equals $1/s$. From TL-Diff, then, $[-1]^N \cdot \mathcal{L}(t^N)(s)$ equals $[\frac{d}{ds}]^N(\widehat{1}(s)) = [\frac{d}{ds}]^N(\frac{1}{s})$. Thus

$$6a: \quad \mathcal{L}(t^N)(s) = \frac{N!}{s^{N+1}}.$$

From LT-Exp, note,

$$*: \quad [\mathcal{L}(e^{Bt})](s) = \frac{1}{s - B}.$$

Recall $\cos(Ft) = \frac{1}{2}[e^{iFt} + e^{-iFt}]$. So $(*)$ gives

$$2 \cdot \mathcal{L}(\cos(Ft))(s) = \frac{1}{s - iF} + \frac{1}{s + iF} = \frac{2s}{s^2 + F^2}.$$

Dividing both sides by 2,

$$\mathcal{L}(\cos(Ft))(s) = \frac{s}{s^2 + F^2}, \quad \text{and similarly}$$

$$6b: \quad \mathcal{L}(\sin Ft)(s) = \frac{F}{s^2 + F^2}.$$

As $\sin()$ is bounded, it has a Laplace transform on \mathbb{R}_+ . Since $\widehat{\sin}(u) = \frac{1}{u^2+1} = \arctan'(u)$, it follows that

$$\int_0^\infty \widehat{\sin}(u) du = \arctan(\infty) - \arctan(0) = \frac{\pi}{2} - 0.$$

Skipping the effort it takes to *justify* applying (5e) at $s=0$, we get that

$$\int_0^\infty \frac{\sin(t)}{t} dt = \mathcal{L}\left(\frac{\sin(t)}{t}\right)(0) \stackrel{\text{by (5e)}}{=} \int_0^\infty \widehat{\sin}.$$

This, together with the previous line, gives

$$6c: \quad \int_0^\infty \frac{\sin(t)}{t} dt = \frac{\pi}{2}.$$

Preliminaries

The **Heaviside fnc** $\mathbf{H}:\mathbb{R}\rightarrow\{0,1\}$ is 0 on $(-\infty,0)$ and is 1 on $[0,+\infty)$. Thus^{♥1}

$$\sum_{K=1}^{\infty} \mathbf{H}(x - K) = \lfloor x \rfloor, \quad \text{for each } x \geq 0.$$

Let δ denote the **Dirac delta “function”**.^{♥2} Write δ_5 for its translate $x \mapsto \delta(x - 5)$. So for P a posreal, and f continuous at P :

$$\int_0^{\infty} f \cdot \delta_P = f(P).$$

In particular, for each f which is continuous:

$$7a: \quad [f \circledast \delta_5](t) = \begin{cases} 0 & , \text{if } t \in [0, 5) \\ f(t - 5) & , \text{if } t \in [5, \infty) \end{cases} .$$

I.e. $f \circledast \delta_5 = \mathbf{T}_5(\mathbf{H} \cdot f)$.

Periodicity. A posreal number P is a “**period** of fnc f ” if

$$\forall x : \quad f(x + P) = f(x).$$

Typically, there is a *smallest* such period, which is called the “**least-period** of f ”.

For a posreal P , let $f_{\langle P \rangle} := f \cdot \mathbf{1}_{[0,P)}$ abbreviate what I will call “ f clipped at P ”. That is, we restrict f to interval $[0, P)$. Thus

$$\widehat{f_{\langle P \rangle}}(s) = \int_0^P e^{-st} f(t) dt.$$

7b: Periodicity Theorem. Suppose $f, g \in \text{Ord}(\mu)$ and P is a posreal. Then

$$1: \mathcal{L}(\mathbf{H}(t - P) \cdot f(t - P))(s) = e^{-sP} \cdot \widehat{f}(s).$$

$$2: \mathcal{L}(\mathbf{H}(t - P)g(t))(s) = e^{-sP} \cdot \mathcal{L}(g(t + P))(s).$$

^{♥1}For a real x , the expression $\lfloor x \rfloor$ is called the **floor** of x ; it is the largest integer less-equal x . So $\lfloor \pi \rfloor$ is 3.

Use DE to mean ‘*differential eqn*’.

^{♥2}It is actually a *Schwartzian distribution*, named after Laurant Schwartz. As distributions, $\mathbf{H}' = \delta$.

3: Suppose now that P is a period of f . Then

$$\widehat{f}(s) = \widehat{f_{\langle P \rangle}}(s) / [1 - e^{-sP}]. \quad \diamond$$

Proof of (7b.1). The LhS equals $\int_P^{\infty} e^{-st} f(t - P) dt$. CoV $x = t - P$ gives $dx = dt$, and thus

$$\text{LhS}(7b.1) = \int_0^{\infty} e^{-s[x+P]} f(x) dx \stackrel{\text{note}}{=} \text{RhS}(7b.1) \quad \blacklozenge$$

Pf (7b.3). Because our “clipped” $f_{\langle P \rangle}(t)$ equals $[\mathbf{H}(t) - \mathbf{H}(t - P)] \cdot f(t)$, it follows that

$$\begin{aligned} \mathcal{L}(f_{\langle P \rangle}) &= \mathcal{L}(\mathbf{H} \cdot f) - \mathcal{L}(\mathbf{H}(t - P) \cdot f(t)) \\ * : \quad \underline{\underline{P \text{ an } f\text{-period}}} \mathcal{L}(\mathbf{H} \cdot f) &= \mathcal{L}(\mathbf{H}(t - P) \cdot f(t - P)). \end{aligned}$$

Our (7b.1) says that $\mathcal{L}(\mathbf{H}(t - P) \cdot f(t - P))(s)$ equals $e^{-sP} \cdot \widehat{f}(s)$. And, always, $\mathcal{L}(\mathbf{H} \cdot f) = \mathcal{L}(f)$, since the Laplace integral is over $[0, \infty)$. So we can re-write (*) as $\widehat{f_{\langle P \rangle}}(s) = [1 - e^{-sP}] \cdot \widehat{f}(s)$. \blacklozenge

Using periodicity. Function $\sin()$ is periodic, with period 2π . So for each posint K :

$$[\sin \otimes \delta_{2\pi K}](t) = \begin{cases} 0 & , \text{if } t \in [0, 2\pi K) \\ \sin(t) & , \text{if } t \in [2\pi K, \infty) \end{cases} .$$

For $N \in \{1, 2, 3, \dots\} \cup \{\infty\}$, define the sum

$$R_N := \sum_{K=1}^N \delta_{2\pi K} .$$

For N finite, then,

$$[\sin \otimes R_N](t) = \lfloor \frac{t}{2\pi} \rfloor \cdot \sin(t)$$

holds when $0 \leq t < 2\pi[N+1]$. There are no convergence problems as $N \nearrow \infty$. So $\forall t \in [0 .. \infty)$:

$$7c: \quad [\sin \otimes R_\infty](t) = \lfloor \frac{t}{2\pi} \rfloor \cdot \sin(t) .$$

Moreover, for each finite N :

$$7c_N: \quad [\sin \otimes R_N](t) = \text{Min}\left(N, \lfloor \frac{t}{2\pi} \rfloor\right) \cdot \sin(t) .$$

Hammering an undamped spring

(One of the textbooks called this *Soldiers marching in cadence*, but the interpretation is less clear, since soldiers already on the bridge are still marching.) With y the unknown-fnc, let us examine DE

$$\mathfrak{D}: \quad y'' + y = \sum_{K=1}^{\infty} \delta_{2\pi K} .$$

Given two complex numbers α, β , let $(\mathfrak{D}_{\alpha, \beta})$ mean the DE together with initial conditions $y(0) = \alpha$ and $y'(0) = \beta$.

Consider the corresponding ZeroTar^{♥3} IVP

$$z'' + z = 0, \quad \text{with } z(0) = \alpha \text{ and } z'(0) = \beta .$$

Its solution is

$$7d: \quad z = [\alpha \cdot \cos] + [\beta \cdot \sin] .$$

^{♥3}Some textbooks call this (yuck!) “homogeneous”.

So the soln, y , to $(\mathfrak{D}_{0,0})$ will, when added to z , give the solution to $(\mathfrak{D}_{\alpha, \beta})$.

Computing,

$$\begin{aligned} \mathcal{L}(y'' + y)(s) &= [s^2 + 1] \cdot \widehat{y}(s) - [s \cdot y(0) + y'(0)] \\ &= [s^2 + 1] \cdot \widehat{y}(s) . \end{aligned}$$

Now (\mathfrak{D}) says that $\mathcal{L}(y'' + y) = \mathcal{L}(R_\infty)$. Dividing,

$$\widehat{y}(s) = \frac{1}{s^2+1} \cdot \widehat{R_\infty}(s) .$$

The Rhs is a product, so its inv-lap-xform is a convolution.

$$y = \sin \otimes R_\infty .$$

In other words, the function

$$7e: \quad y_{\alpha, \beta}(t) := \alpha \cos(t) + \beta \sin(t) + \lfloor \frac{t}{2\pi} \rfloor \cdot \sin(t)$$

is the general solution to IVP $(\mathfrak{D}_{\alpha, \beta})$.

At this point, it would be good to have a careful sketch of $y_{\alpha, \beta}()$. One could also analyze, at each time t , the potential and kinetic energy in the spring. If the spring is damped, is there a net absorption, or loss, of energy?

Suppose the tapping on the spring was not aligned with the resonant frequency of the spring—would the spring nonetheless absorb energy?

Random hammering. Times $0 < G_1 < G_2 < \dots$ define

$$R_G := \sum_{K=1}^{\infty} \delta_{G_K} .$$

The above reasoning applied to DE $y'' + y = R_G$ yields $y = \sin \otimes R_G$, as before. Computing the convolution gives

$$7f: \quad y(t) = \sum_{K=1}^{\infty} \mathbf{H}(t - G_K) \cdot \sin(t - G_K) .$$

Adding $[\alpha \cos(t) + \beta \sin(t)]$ hands us the general solution $y_{\alpha, \beta}$.

Square Wave

For a period $P > 0$ and a **duty cycle** $U \in [0, 1]$, define the **square-wave fnc** $\text{SW}_{P,U}()$ to be the P -periodic fnc mapping $\mathbb{R} \rightarrow \mathbb{R}$, which agrees with $\mathbf{1}_{[0,UP)}$ on interval $[0, P)$; the **pulse width** is UP .

So $\text{SW}_{P,0} \equiv 0$ and $\text{SW}_{P,1} \equiv 1$. Thus we expect

$$*: \quad \widehat{\text{SW}}_{P,0} \equiv 0 \quad \text{and} \quad \widehat{\text{SW}}_{P,1}(s) = 1/s.$$

Setting $f := \text{SW}_{P,U}$, our Periodicity Thm says that $\widehat{f_{(P)}}(s)$ equals

$$\begin{aligned} \int_0^P e^{-st} \cdot \mathbf{1}_{[0,UP)}(t) dt &\stackrel{\text{since } U \leq 1}{=} \int_0^{UP} e^{-st} dt \\ &= \frac{1}{s} \cdot [1 - e^{-s \cdot UP}]. \end{aligned}$$

Consequently,

$$\dagger 1: \quad \widehat{\text{SW}}_{P,U}(s) = \frac{1}{s} \cdot M_{P,U}(s),$$

where this multiplier fnc $M_{P,U}$ is

$$\dagger 2: \quad M_{P,U}(s) = \frac{1 - e^{-UP \cdot s}}{1 - e^{-P \cdot s}} \stackrel{\text{note}}{=} \frac{e^{P \cdot s} - e^{FP \cdot s}}{e^{P \cdot s} - 1};$$

here, $F := 1 - U$ represents the “off” part of the cycle. Happily, $(\dagger 1, \dagger 2)$ is consistent with $(*)$.

Square-Wave into Spring. Laplace transforming $y'' + y = \text{SW}_{P,U}$ yields

$$\text{Y}: \quad [s^2 + 1] \cdot \widehat{y}(s) = \widehat{\text{SW}}_{P,U}(s).$$

As before, then,

$$8: \quad y = \sin \circledast \text{SW}_{P,U}.$$

Alternatively, write (Y) as

$$\begin{aligned} [s^2 + 1] \cdot \widehat{y}(s) &= \frac{1}{s} \cdot M_{P,U}(s). \quad \text{So} \\ \widehat{y}(s) &= \frac{1}{s^2 + 1} \cdot \frac{1}{s} \cdot M_{P,U}(s). \quad \text{Thus} \end{aligned}$$

$$\begin{aligned} 8': \quad y &= \sin \circledast \mathbf{1} \circledast \mathcal{L}^{-1}(M_{P,U}) \\ &= [1 - \cos] \circledast \mathcal{L}^{-1}(M_{P,U}). \end{aligned}$$