

## γ and Γ()

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### EM-number

Recall that the *harmonic numbers*  $\mathcal{H}_N := \sum_{k=1}^N \frac{1}{k}$  upper/lower bound  $\log()$ , in that

$$\mathcal{H}_{N-1} > \log(N) > \mathcal{H}_N - 1.$$

The *Euler-Mascheroni* number<sup>♥1</sup> is defined as the asymptotic discrepancy,

$$\begin{aligned} \gamma &:= \lim_{N \nearrow \infty} [\mathcal{H}_{N-1} - \log(N)] \\ \text{1:} \quad &= \lim_{N \nearrow \infty} [\mathcal{H}_N - \log(N)] = \int_1^\infty \left[ \frac{1}{\lfloor t \rfloor} - \frac{1}{t} \right] dt. \end{aligned}$$

Although estimate  $\gamma \approx 0.578^-$  does not help with rationality, it is reassuring to know.

**1a: Lemma.** *The  $N^{\text{th}}$  harmonic number satisfies*

$$\text{*1:} \quad \mathcal{H}_N = \sum_{k=1}^{\infty} \frac{N}{k[k+N]}.$$

Also

$$\text{*2:} \quad \frac{N}{k[k+N]} = \int_0^1 \frac{t^{k/N}}{k} dt.$$

Consequently,

$$\text{*3:} \quad \mathcal{H}_N = -\int_0^1 \log(1 - t^{1/N}) dt. \quad \diamond$$

**Pf (\*1).** Sum  $S_\ell := \sum_{k=1}^{\ell} \frac{N}{k[k+N]} = \sum_{k=1}^{\ell} \left[ \frac{1}{k} - \frac{1}{k+N} \right]$  telescopes. Once  $\ell \geq N$ , then,

$$S_\ell = \left[ \sum_{k=1}^N \frac{1}{k} \right] - \left[ \sum_{k=1}^N \frac{1}{\ell+k} \right].$$

The righthand-sum goes to zero, as  $\ell \nearrow \infty$ . ♦

**Pf of (\*2).** Observe  $\int t^{\frac{k}{N}} dt = \frac{N}{k+N} t^{\frac{k+N}{N}}$ . Hence RhS(\*2) equals

$$\frac{N}{k[k+N]} \cdot t^{\frac{k+N}{N}} \Big|_{t=0}^{t=1} = \text{LhS}(*2). \quad \diamond$$

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<sup>♥1</sup>As of Nov.2017, it is unknown if the *EM-number*  $\gamma$  is rational or irrational.

**Pf of (\*3).** One Taylor-series expansion for  $\log$  is

$$-\log(1-x) \stackrel{\text{for } |x| < 1}{=} \sum_{k=1}^{\infty} \frac{x^k}{k}.$$

For  $0 < t < 1$ , then,

$$-\log(1 - t^{1/N}) = \sum_{k=1}^{\infty} \frac{t^{\frac{k}{N}}}{k}.$$

It is valid to commute  $\int_0^1$  with the sum, as all the summands have the same sign. Thus

$$\begin{aligned} -\int_0^1 \log(1 - t^{1/N}) dt &= \sum_{k=1}^{\infty} \int_0^1 \frac{t^{\frac{k}{N}}}{k} dt \\ &= \sum_{k=1}^{\infty} \frac{N}{k[k+N]} = \mathcal{H}_N. \quad \diamond \end{aligned}$$

**1b: Lemma.** *For all  $t$  with  $0 < t < 1$ ,*

$$\dagger: \quad \lim_{N \nearrow \infty} N \cdot [1 - t^{1/N}] = -\log(t).$$

$$\ddagger: \quad \lim_{N \nearrow \infty} \log(N \cdot [1 - t^{1/N}]) = \log(-\log(t)). \quad \diamond$$

**Proof.** Setting  $h := \frac{1}{N}$ , we can rewrite limit (†) as

$$-\lim_{h \searrow 0} \frac{t^h - 1}{h} \stackrel{\text{note}}{=} -f'(0),$$

where  $f(h) := t^h$ . By definition,  $f(h) = \exp(\log(t) \cdot h)$ . So  $f'(h) = \log(t) \cdot f(h)$ . Thus  $f'(0) = \log(t)$ .

Lastly, (‡) holds, since [the outer]  $\log$  is continuous. ♦

**1c: γ-Γ Thm.** *Using fnc Γ() from the next section,*

$$\gamma = \int_0^1 \log(-\log(t)) dt = -\Gamma'(1). \quad \diamond$$

**Proof.** Eqns (1) and (\*3) say  $\gamma$  is the limit of

$$\begin{aligned} \mathcal{H}_N - \log(N) &= \left[ -\int_0^1 \log(1 - t^{1/N}) dt \right] - \int_0^1 \log(N) dt \\ &= -\int_0^1 \left[ \log(1 - t^{1/N}) + \log(N) \right] dt \\ &= -\int_0^1 \log(N \cdot [1 - t^{1/N}]) dt. \end{aligned}$$

Skipping the justification needed to pass  $\lim_{N \nearrow \infty}$  through the integral sign,

$$\begin{aligned} \gamma &= -\int_0^1 \left[ \lim_{N \nearrow \infty} \log(N \cdot [1 - t^{1/N}]) \right] dt \\ &\stackrel{\text{by } (\ddagger)}{=} -\int_0^1 \log(-\log(t)) dt. \end{aligned}$$

**Exponential CoV.** To compute this last integral, we use CoV  $u = -\log(t)$ . I.e,  $t = e^{-u}$ , so  $\frac{dt}{du} = -e^{-u}$  and thus  $dt = -e^{-u} du$ .

As  $0 \nearrow t \nearrow 1$ , remark that  $\infty \searrow u \searrow 0$ . Hence,

$$\int_0^1 \log(-\log(t)) dt \stackrel{\text{CoV}}{=} \int_\infty^0 \log(u) \overbrace{[-e^{-u}]}^{dt} du = \int_0^\infty \log(u) e^{-u} du.$$

And this last equals  $\Gamma'(1)$ , courtesy (2d) ♦

**Gamma function**

The **Gamma fnc** arises in volumes of  $N$ -dimensional balls, and in Laplace transforms. Leaving motivation for later, define the Gamma fnc by

2:  $\Gamma(z) := \int_0^\infty t^{z-1} e^{-t} dt, \quad \text{for } z \in \mathbb{C} \text{ with } \text{Re}(z) > 0,$

When  $\text{Re}(w) > 0$ , note  $\lim_{t \searrow 0} [t^w e^{-t}]$  is zero. Thus

$$t^w [-e^{-t}] \Big|_{t=0}^{t=\infty} = 0.$$

For  $\text{Re}(z) > 1$ , then, integration by parts produces

$$\Gamma(z) = t^{z-1} \cdot [-e^{-t}] \Big|_{t=0}^{t=\infty} - \int_0^\infty [z-1] t^{z-2} \cdot [-e^{-t}] dt. \text{ So}$$

2a:  $\Gamma(z) = [z-1] \cdot \Gamma(z-1), \quad \text{for } \text{Re}(z) > 1.$

**Analytic continuation.** Writing  $w := z-1$ ,

2b:  $\Gamma(w) = \frac{1}{w} \cdot \Gamma(w+1)$

from (2a). As (2d) will show, our  $\Gamma()$  is  $\infty$ ly differentiable, hence is *analytic* courtesy the Cauchy-Goursat theorem. Use (2b) to iteratively extend  $\Gamma()$  to the complex plane. This  $\Gamma()$  is **analytic on  $\mathbb{C} \setminus \{0, -1, -2, -3, \dots\}$** . Further,  $\Gamma()$  is **meromorphic on  $\mathbb{C}$** , with **simple poles at  $0, -1, -2, \dots$** .

To see that  $\Gamma()$  has a simple pole at 0, it suffices to note that  $\Gamma(1) = \int_0^\infty e^{-t} dt = 1$  is non-zero.

Since  $\Gamma(1) = 1$ , our (2a) implies

2c:  $\Gamma(n) = [n-1]!, \quad \text{for } n \in \mathbb{Z}_+.$

As a consequence, binomial/multinomial coefficients can be generalized using  $\Gamma$ , and many of the identities extend.

**Calculus.** Differentiating under integral sign is valid in (2). For  $k = 0, 1, \dots$ , applying  $\frac{d}{dz}$  gives

$\Gamma^{(k)}(z) = \int_0^\infty t^{z-1} e^{-t} [\log(t)]^k dt, \text{ when } \text{Re}(z) > 0.$   
 2d: In particular,  $\Gamma'(1) = \int_0^\infty e^{-t} \log(t) dt \stackrel{\text{by (1c)}}{=} -\gamma.$

Use  $\hat{f}$  to mean the **Laplace transform** of  $f$ , where

$$\hat{f}(s) = [\mathcal{L}(f)](s) := \int_0^\infty e^{-st} \cdot f(t) \cdot dt.$$

2e: **Theorem.** *The Laplace transform of logarithm, for  $s > 0$ , is*

$$\widehat{\log}(s) = \frac{1}{s} [\Gamma'(1) - \log(s)]$$

$$\dagger: = \frac{-1}{s} [\gamma + \log(s)].$$

Fix  $z$  with  $\text{Re}(z) > -1$ . For  $t > 0$ , define  $P(t) := t^z$ . Then

‡:  $\hat{P}(s) = \frac{\Gamma(z+1)}{s^{z+1}}, \quad \text{for } s > 0. \quad \diamond$

**Pf of (†).** Fix an  $s > 0$ . With  $u := st$ , then,  $dt = \frac{1}{s} du$ . As  $0 \nearrow t \nearrow \infty$ , our  $0 \nearrow u \nearrow \infty$ , since  $s > 0$ . Thus  $\widehat{\log}(s)$  equals

$$\int_0^\infty e^{-st} \log(t) \cdot dt \stackrel{\text{CoV}}{=} \int_0^\infty e^{-u} \overbrace{\log\left(\frac{u}{s}\right)}^{\log(t)} \cdot \overbrace{\frac{1}{s}}^{dt} du. \text{ So,}$$

$$s \cdot \widehat{\log}(s) = \int_0^\infty e^{-u} [\log(u) - \log(s)] du$$

$$= \Gamma'(1) - \log(s)\Gamma(1) = \Gamma'(1) - \log(s),$$

as claimed. ♦

**Pf of (‡).** With  $u := st$ , as before  $dt = \frac{1}{s} du$ . So  $\hat{P}(s)$  equals

$$\int_0^\infty e^{-st} P(t) \cdot dt \stackrel{\text{since } s > 0}{=} \int_0^\infty e^{-u} \left[\frac{u}{s}\right]^z \cdot \overbrace{\frac{1}{s}}^{dt} du$$

$$= \frac{1}{s^{z+1}} \int_0^\infty e^{-u} u^z du. \quad \diamond$$

As our last fact in this introduction to  $\Gamma()$ , let's show that

2f:  $\Gamma\left(\frac{1}{2}\right) = \sqrt{\pi}$ .

We employ substitution  $t = x^2$ . Thus  $dt = 2x dx$ . As  $0 \nearrow t \nearrow \infty$ , note  $0 \nearrow x \nearrow \infty$ . Recall  $\Gamma\left(\frac{1}{2}\right)$  equals

$$\begin{aligned} \int_0^\infty e^{-t} t^{\frac{1}{2}-1} dt &= \int_0^\infty e^{-t} \frac{1}{t^{1/2}} \cdot dt \\ &\stackrel{\text{CoV}}{=} \int_0^\infty e^{-x^2} \frac{1}{x} \cdot 2x dx \\ &= 2 \int_0^\infty e^{-x^2} dx = \int_{-\infty}^\infty e^{-x^2} dx. \end{aligned}$$

This last integral equals  $\sqrt{\pi}$ , as shown in the next section by means of the famous Polar-Coordinate Trick.

### Polar-coordinate Trick

Let  $J := \int_{-\infty}^{+\infty} e^{-x^2} dx$ . We use the *PCT* (“polar coordinate trick”) to show that  $J = \sqrt{\pi}$ . We integrate the cartesian-square of the integrand to conclude that

3.1: 
$$\begin{aligned} J^2 &= \left[ \int_{-\infty}^{+\infty} e^{-[x^2]} dx \right] \cdot \left[ \int_{-\infty}^{+\infty} e^{-[y^2]} dy \right] \\ &= \int_{-\infty}^{+\infty} e^{-[x^2+y^2]} \cdot d(x, y) \\ &= \int_0^{+\infty} e^{-r^2} \cdot \underbrace{2\pi r \cdot dr}_{\substack{\text{Area of radius-}r \text{ annulus} \\ \text{of thickness } dr}}. \end{aligned}$$

Hence  $J^2 = \pi \cdot [-e^{-r^2}]|_{r=0}^{r=+\infty} = \pi$ . Since  $J$  is the integral of a non-negative fnc, nec.  $J \geq 0$ . Thus

3.2: 
$$\int_{-\infty}^{+\infty} e^{-x^2} dx = \sqrt{\pi}.$$

## Further $\Gamma$ results

This section will always be “in progress”.

It follows from the defn of  $\Gamma$  that  $\Gamma(\bar{z}) = \overline{\Gamma(z)}$ .

4: Prop. For  $N$  a natnum,

$$\Gamma'(N+1) = N! \cdot [\mathcal{H}_N - \gamma].$$
 ◇

*Proof.* Set  $g(N+1) := N! \cdot [\mathcal{H}_N - \gamma]$ . For  $N \geq 1$ , note,  $g(N+1) - N \cdot g(N)$  equals

$$N! \cdot [\mathcal{H}_N - \gamma] - N! \cdot [\mathcal{H}_{N-1} - \gamma] = N! \cdot [\mathcal{H}_N - \mathcal{H}_{N-1}].$$

Restated,

†: 
$$g(N+1) - N \cdot g(N) = [N-1]!$$

Recall  $\Gamma(w+1) = w \cdot \Gamma(w)$  from (2b). Differentiating,

$$\Gamma'(w+1) = 1 \cdot \Gamma(w) + w \Gamma'(w).$$

Hence

‡: 
$$\Gamma'(N+1) - N \cdot \Gamma'(N) = \Gamma(N) \stackrel{\text{by (2c)}}{=} [N-1]!$$

So  $\Gamma'$  and  $g$  satisfy the same recurrence relation. And

$$\Gamma'(1) \stackrel{\text{by (2c)}}{=} -\gamma \stackrel{\text{def}}{=} g(0+1)$$

handles the base case. ◇

5: Euler's reflection formula. For all  $z \in \mathbb{C}$ :

$$\Gamma(z) \cdot \Gamma(1-z) = \frac{\pi}{\sin(\pi z)}.$$
 ◇

*Proof.* Challenging... [Courtesy analytic-continuation, IT-SProve (5) for a seq. of  $z$  converging to a point of analyticity. We will prove (5) for real  $0 < Z < 1$ .] ● ● ● ◇

6: Lemma. For real  $Z$  with  $0 < Z < 1$ :

$$\Gamma(z) \cdot \Gamma(1-z) = \int_0^\infty \frac{\rho^{Z-1}}{\rho+1} d\rho.$$
 ◇

*Proof.* Difference  $1-Z$  is in  $(0, 1)$ , so  $\Gamma(Z) \cdot \Gamma(1-Z)$  equals

$$\begin{aligned} \mathbf{I} &:= \left[ \int_0^\infty x^{Z-1} e^{-x} dx \right] \int_0^\infty t^{[1-Z]-1} e^{-t} dt \\ &= \int_0^\infty \int_0^\infty e^{-[x+t]} x^{Z-1} t^{-Z} dt dx. \end{aligned}$$

We use CoV [sum]  $\sigma = t+x$ , and [ratio]  $\rho = \frac{x}{t}$ . For a fixed positive sum  $\sigma = t+x$ , note that  $x$  can be made arbitrarily close to zero, or to  $\sigma$ . Hence the ratio  $\rho = \frac{x}{t}$  varies over  $\mathbb{R}_+$ . Thus as  $(t, x)$  varies over  $\mathbb{R}_+ \times \mathbb{R}_+$ , so does  $(\sigma, \rho)$ ; i.e, after the CoV, the limits-of-integration will be  $\int_0^\infty \int_0^\infty$ .

**Jacobian.** The Jacobian matrix is

$$\begin{bmatrix} \partial\sigma/\partial t & \partial\rho/\partial t \\ \partial\sigma/\partial x & \partial\rho/\partial x \end{bmatrix} = \begin{bmatrix} 1 & -x/t^2 \\ 1 & 1/t \end{bmatrix}. \quad \begin{array}{l} \text{Its determinant} \\ \text{is thus} \end{array}$$

$$\frac{\partial(\sigma, \rho)}{\partial(t, x)} := \frac{1}{t} - \frac{-x}{t^2} \stackrel{\text{note}}{=} \frac{\sigma}{t^2}.$$

Viewing  $\sigma$  and  $\rho$  as functions of  $t$  and  $x$ , we rewrite **I** as

$$*: \int_0^\infty \int_0^\infty e^{-[t+x]} \cdot x^{Z-1} t^{-Z} \frac{t^2}{\sigma} \cdot \underbrace{\frac{\partial(\sigma, \rho)}{\partial(t, x)} dt dx}$$

Now  $\sigma = x + t = \rho t + t = [\rho + 1]t$ , so  $\frac{t}{\sigma} = \frac{1}{\rho+1}$ . And  $x^{Z-1} t^{-Z} \frac{t^2}{\sigma} = [\rho t]^{Z-1} t^{-Z} \frac{t^2}{\sigma} = \rho^{Z-1} \frac{t}{\sigma} = \frac{\rho^{Z-1}}{\rho+1}$ .

Our CoV applied to (\*) says  $\Gamma(Z) \cdot \Gamma(1-Z)$  equals

$$7: \int_0^\infty \int_0^\infty e^{-\sigma} \cdot \frac{\rho^{Z-1}}{\rho+1} \cdot d\sigma d\rho = \int_0^\infty \frac{\rho^{Z-1}}{\rho+1} d\rho,$$

since  $\int_0^\infty e^{-\sigma} d\sigma$  equals 1. ♦